

Symbolic methods to construct a cusp, breathers, kink, rogue waves and some soliton waves solutions of nonlinear partial differential equations

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Abstract A cusp, bright breathers, dark breathers, kink, bright rogue waves and some soliton waves solutions are obtained by using the $\exp(-\phi(\xi))$ -expansion method for the fourth order Benjamin-Ono equation and BBM equations. The obtained solutions might be indicated and meaningful for narrating the physical phenomena in the real-world. For compatible values of the arbitrary parameter included in the solution, we plot the 3D surface of the all obtained solutions which are shown in Figure 1 to Figure 10.

Keywords. The $\exp(-\phi(\xi))$ -expansion method, The fourth order Benjamin-Ono equation, BBM equation, Traveling wave solutions, Nonlinear evolution equation.

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1. INTRODUCTION

Finding the exact and analytical solutions to nonlinear partial differential equations (NLPDEs) play fundamental and imperative problems in mathematical physics, applied sciences and engineering. As a result, various groups of physicist and mathematicians have been working tirelessly to improve useful methods for providing different types soliton form solutions to NLPDEs. For this reason, several methods have been developed to search the analytical solutions, such as the novel (G'/G) -expansion method [2, 3, 4, 6, 8, 9, 10], the extended exp-function method [26, 28, 29], the modified exp-function method [27, 30], the sine-Gordon expansion method [13, 14], the homogeneous balance method [16, 25], the simplest equation method [23], the inverse scattering transform method [1, 33], the Hirota's bilinear method [38], the tanh-function method [37], Bell-polynomial method [41], the extended tanh-function method [17, 34, 40], the Exp-function method [20], the sine-cosine method [36], the

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Backlund transformation method [22], the modified Exp-function method [32], the Jacobi elliptic function expansion method [21], the $\exp(-\phi(\xi))$ -expansion method [5, 7] and so on.

The objective of this article is to derive some fresh and wide-applicable a cusp, bright breathers, dark breathers, kink, bright rogue waves and some soliton waves solutions to the fourth order Benjamin-Ono equation and BBM equations through the $\exp(-\phi(\xi))$ -expansion method which are discussed in section 3. The new analytical solutions performed in this paper are characterized by exponent, trigonometric, rational, and hyperbolic functions. However, we observe that a cusp, bright bright breathers, dark breathers, kink, bright rogue waves and some soliton waves solutions are obtained to the fourth order Benjamin-Ono equation and BBM equations have not been reported previously.

The synopsis of this paper as follows: In Section 2, we give the algorithm of the $\exp(-\phi(\xi))$ -expansion method. In Section 3, new solutions of the fourth order Benjamin-Ono equation and BBM equations are formulated through the $\exp(-\phi(\xi))$ -expansion method. In Section 4, graphical representations and numerical experiments of the derived solutions are depicted. Finally, the conclusion of our study is given.

2. ALGORITHM OF $\exp(-\phi(\xi))$ -EXPANSION METHOD

The following is given as the general nonlinear PDEs (the fourth order Benjamin-Ono equation and BBM equations) with two variables x and t as

$$P(v, v_t, v_x, v_{tt}, v_{xx}, v_{xt}, \dots) = 0, \quad (2.1)$$

where $v(x, t)$ is an unknown function and P is a polynomial in $v(x, t)$.

Step 1: The traveling wave variable

$$v(x, t) = v(\xi), \xi = x + Vt, \quad (2.2)$$

where V is the speed of the traveling wave and the traveling wave variable (2.2) converts equation (2.1) into the ordinary differential equation

$$R(v, -Vv', v', v', v', v'', v'', v'', V^2v'', \dots) = 0 \quad (2.3)$$

where R is a polynomial in v and its derivatives and the superscripts indicate the ordinary derivatives with respect to ξ .

Step 2: Suppose that the traveling wave solution of equation (2.3) can be expressed as

$$v(\xi) = \sum_{i=0}^N A_i (\exp(-\phi(\xi)))^i, \quad (2.4)$$

where $A_i (i = 0, \dots, n)$ are constants to be determined, such that $A_N \neq 0$ and $\phi = \phi(\xi)$ satisfies the following ordinary differential equation:

$$(\phi(\xi))' = \exp(-\phi(\xi)) + \mu \exp(\phi(\xi)) + \lambda, \quad (2.5)$$

where $A_N, \dots, V, \lambda, \mu$ are constants to be determined latter.

Step 3: For a given ansatz equation (for example, the ansatz equation is $(\phi(\xi))' =$



$exp(-\phi(\xi)) + \mu exp(\phi(\xi)) + \lambda$ in this paper), the form of v is decided and the homogeneous balance method is used on equation (2.3) to find the coefficients of v .

Step 4: The homogeneous balance method is used to solve the ansatz equation.

Step 5: Finally, the solitary wave solutions of equation (2.1) are obtained by combining steps 3 and 4.

Next, we have five solutions including trigonometric, hyperbolic, exponent and rational function structure of equation (2.5).

When $\mu \neq 0, \lambda^2 - 4\mu > 0,$, the solution of equation (2.5) is

$$(\phi(\xi)) = \ln\left(\frac{-\sqrt{\lambda^2 - 4\mu} \tanh\left(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\right)(\xi + E) - \lambda}{2\mu}\right). \tag{2.6}$$

When $\mu \neq 0, \lambda^2 - 4\mu < 0,$ the solution of equation (2.5) is

$$(\phi(\xi)) = \ln\left(\frac{\sqrt{4\mu - \lambda^2} \tan\left(\frac{\sqrt{4\mu - \lambda^2}}{2}\right)(\xi + E) - \lambda}{2\mu}\right). \tag{2.7}$$

When $\mu = 0, \lambda \neq 0, \lambda^2 - 4\mu > 0,$, the solution of equation (2.5) is

$$(\phi(\xi)) = \ln\left(\frac{\lambda}{exp(\lambda(\xi + E)) - 1}\right). \tag{2.8}$$

When $\mu \neq 0, \lambda \neq 0, \lambda^2 - 4\mu = 0,$, the solution of equation (2.5) is

$$(\phi(\xi)) = \ln\left(-\frac{(2\lambda(\xi + E) + 2)}{\lambda^2(\xi + E)}\right). \tag{2.9}$$

When $\mu = 0, \lambda = 0, \lambda^2 - 4\mu = 0,$, the solution of equation (2.5) is

$$(\phi(\xi)) = \ln(\xi + E). \tag{2.10}$$

3. FORMULATION OF SOLUTIONS

In this section, the method is used to construct a cusp, bright breathers, dark breathers, kink, bright rogue waves and some soliton waves solutions for the fourth order Benjamin-Ono equation and BBM equations which are very important nonlinear evolution equations in the field of nonlinear dynamics.

3.1. The fourth order Benjamin-Ono equation. In this subsection, we will present the $exp(-\phi(\xi))$ -expansion method to construct bright breathers, dark breathers, bright rogue waves and multiple soliton waves solutions of the fourth order Benjamin-Ono equation which are illustrated in Figure 1 to Figure 5. The fourth order Benjamin-Ono equation is of the form,

$$v_{tt} + \alpha(v^2)_{xx} + \beta v_{xxxx} = 0, \tag{3.1}$$

where α and β are real non-zero parameters and $(v^2)_{xx}$ is the nonlinear term and v_{xxxx} is dispersion term. In mathematics, the fourth order Benjamin-Ono equation is a significant nonlinear partial integro differential equation that derives one-dimensional internal waves in deep water and it plays a important problem in many scientific applications, such as fluid dynamics, nonlinear optics, plasma physics [18, 35, 19, 15, 39].



By using traveling wave variable $v(\xi) = v(x, t)$, $\xi = x - Vt$, Equation (3.1) converts into a nonlinear ordinary differential equation

$$V^2 v'' + \alpha(v^2)'' + \beta v^{iv} = 0. \quad (3.2)$$

Performing the first and second integrations of equation (3.2), we get

$$K + V^2 v + \alpha(v^2) + \beta v'' = 0, \quad (3.3)$$

where K is constant. Making the homogeneous balance between v^2 and v'' , we get $N = 2$. Putting the value of N in equation (3.3) our solution is of the form

$$v(\xi) = A_0 + A_1(\exp(-\phi(\xi))) + A_2(\exp(-\phi(\xi))), \quad (3.4)$$

where the coefficients A_0 , A_1 and A_2 are constants to be evaluated.

Putting the values of equation (3.4) into equation (3.3) and then equating each coefficients of $\exp(-\phi(\xi))$ to zero, we get

$$6\beta A_2 + \alpha A_2^2 = 0, \quad (3.5)$$

$$10\beta A_2 \lambda + 2\alpha A_1 A_2 + 2\beta A_1 = 0, \quad (3.6)$$

$$V^2 A_2 + \alpha A_1^2 + 3\beta A_1 \lambda + 2\alpha A_0 A_2 + 8\beta A_2 \mu + 4\beta A_2 \lambda^2 = 0, \quad (3.7)$$

$$V^2 A_1 + 2\alpha A_1 \mu + 6\beta A_2 \mu \lambda + \beta A_1 \lambda^2 + 2\alpha A_0 A_1 = 0, \quad (3.8)$$

$$V^2 A_0 + 2\alpha A_0^2 + 2\beta A_2 \mu^2 + K + \beta A_1 \lambda \mu = 0, \quad (3.9)$$

Using algebraic software Maple, we solve the equation (3.5) to equation (3.9)

$$K = \frac{A_0^2 \alpha^2 + 12\beta^2 \mu^2 + 6\beta^2 \lambda^2 \mu + 8A_0 \alpha \mu \beta + A_0 \alpha \lambda^2 \beta}{\alpha},$$

$$V = \sqrt{-(2\alpha A_0 + 8\beta \mu + \beta \lambda^2)},$$

$$A_0 = A_0, A_1 = -\frac{6\beta \lambda}{\alpha}, A_2 = -\frac{6\beta}{\alpha}.$$

where λ and μ are constants. Now setting the values of V, A_0, A_1, A_2 into equation (3.4), we have

$$v(\xi) = A_0 - \frac{6\beta \lambda}{\alpha}(\exp(-\phi(\xi))) - \frac{6\beta}{\alpha}(\exp(-\phi(\xi)))^2, \quad (3.10)$$

where $\xi = x - (\sqrt{-(2\alpha A_0 + 8\beta \mu + \beta \lambda^2)})t$. A substitution of the equation (2.6) to equation (2.10) into equation (3.10), leads to the following five traveling wave solutions of the fourth order Benjamin-Ono equation.



When $\mu \neq 0, \lambda^2 - 4\mu > 0$, the solution of equation (3.1) is

$$v_1(\xi) = A_0 + \frac{6\beta\lambda}{\alpha} \left(\frac{2\mu}{\sqrt{\lambda^2 - 4\mu} \tanh\left(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\right)(\xi + E) + \lambda} \right) - \frac{6\beta}{\alpha} \left(\frac{2\mu}{\sqrt{\lambda^2 - 4\mu} \tanh\left(\frac{\sqrt{\lambda^2 - 4\mu}}{2}\right)(\xi + E) + \lambda} \right)^2. \tag{3.11}$$

When $\mu \neq 0, \lambda^2 - 4\mu < 0$, the solution of equation (3.1) is

$$v_2(\xi) = A_0 - \frac{6\beta\lambda}{\alpha} \left(\frac{2\mu}{\sqrt{4\mu - \lambda^2} \tan\left(\frac{\sqrt{4\mu - \lambda^2}}{2}\right)(\xi + E) - \lambda} \right) - \frac{6\beta}{\alpha} \left(\frac{2\mu}{\sqrt{4\mu - \lambda^2} \tan\left(\frac{\sqrt{4\mu - \lambda^2}}{2}\right)(\xi + E) - \lambda} \right)^2. \tag{3.12}$$

When $\mu = 0, \lambda \neq 0, \lambda^2 - 4\mu > 0$, the solution of equation (3.1) is

$$v_3(\xi) = A_0 - \frac{6\beta\lambda}{\alpha} \left(\frac{\lambda}{\exp(\lambda(\xi + E)) - 1} \right) - \frac{6\beta}{\alpha} \left(\frac{\lambda}{\exp(\lambda(\xi + E)) - 1} \right)^2. \tag{3.13}$$

When $\mu \neq 0, \lambda \neq 0, \lambda^2 - 4\mu = 0$, the solution of equation (3.1) is

$$v_4(\xi) = A_0 + \frac{6\beta\lambda}{\alpha} \left(\frac{\lambda^2(\xi + E)}{2(\lambda(\xi + E)) + 2} \right) - \frac{6\beta}{\alpha} \left(\frac{\lambda^2(\xi + E)}{2(\lambda(\xi + E)) + 2} \right)^2 \tag{3.14}$$

when $\mu = 0, \lambda = 0, \lambda^2 - 4\mu = 0$, the solution of equation (3.1) is

$$v_5(\xi) = A_0 - \frac{6\beta\lambda}{\alpha} \left(\frac{1}{(\xi + E)} - \frac{6\beta}{\alpha} \frac{1}{(\xi + E)^2} \right) \tag{3.15}$$

3.2. The BBM equation. In this subsection, we will present the $\exp(-\phi(\xi))$ - expansion method to construct a cusp, bright breathers, dark breathers, kink and multiple soliton waves solutions of the Benjamin-Bona-Mahony equation which are shown in Figure 6 to Figure 10. The Benjamin-Bona-Mahony equation is of the form [12, 24, 31]

$$v_t + v_x + avv_x - bv_{xxt} = 0. \tag{3.16}$$

where a and b are constants and $v(x, t)$ is the function of the space variables x for the displacement of the water surface at location in the context of shallow water waves and time variable t . It was established as a model for the unidirectional propagation of small amplitude long waves on the surface of water in a channel in nonlinear dispersive media [11]. According to the shallow water waves, it is applicable to the study of drift waves in the Rossby waves or plasma waves in rotating fluids and also the acoustic-gravity waves in compressible fluids, the acoustic waves in an harmonic crystals and the hydromagnetic waves in cold plasma [11].

By using traveling wave variable $\xi = x - \beta t$ carries equation (3.16) into a nonlinear ordinary differential equation

$$(1 - \beta)v + \frac{a}{2}v^2 - b\beta v'' + K = 0. \tag{3.17}$$



Performing the first and second integrations of equation (3.17) and making the homogeneous balance between v^2 and v''' , we get $N = 2$. Putting the value of N in equation (3.17). Our solution is of the form

$$v(\xi) = A_0 + A_1(\exp(-\phi(\xi))) + A_2(\exp(-\phi(\xi))) \quad (3.18)$$

where the coefficients A_0 , A_1 and A_2 are constants to be evaluated.

Putting the values of equation (3.18) into equation (3.17) and then equating each coefficients of $\exp(-\phi(\xi))$ to zero, we get

$$-6\beta b A_2 + \frac{1}{2} a A_2^2 = 0, \quad (3.19)$$

$$-10\beta b A_2 \lambda + a A_1 A_2 - 2\beta b A_1 = 0, \quad (3.20)$$

$$-\beta A_2 + A_2 + \frac{1}{2} a A_1^2 - 3\beta b A_1 \lambda + a A_0 A_2 - 8\beta b A_2 \mu - 4\beta b A_2 \lambda^2 = 0, \quad (3.21)$$

$$A_1 - \beta A_1 - 2\beta b A_1 \mu - 6\beta b A_2 \mu \lambda - \beta b A_1 \lambda^2 + a A_0 A_1 = 0, \quad (3.22)$$

$$A_0 - \beta b A_1 \lambda \mu - \beta A_0 + K + \frac{1}{2} a A_0^2 - 2\beta b A_2 \mu^2 = 0, \quad (3.23)$$

Using algebraic software Maple, we solve the equation (3.19) to equation (3.23)

$$K = -\frac{2\beta - 1 - 8\beta^2 b^2 \lambda^2 \mu - \beta^2 + \lambda^4 b^2 \beta^2 + 16\mu^2 b^2 \beta^2}{2a},$$

$$\beta = \beta, A_0 = \frac{\lambda^2 b \beta + 8\mu b \beta + \beta - 1}{a},$$

$$A_1 = \frac{12b\beta\lambda}{a}, A_2 = \frac{12b\beta}{a},$$

where λ and μ are constants.

Now setting the values of V, A_0, A_1, A_2 into equation (3.18), we have

$$v(\xi) = \frac{\lambda^2 b \beta + 8\mu b \beta + \beta - 1}{a} + \frac{12b\beta\lambda}{a} (\exp(-\phi(\xi))) + \frac{12b\beta}{a} (\exp(-\phi(\xi)))^2, \quad (3.24)$$

where $\xi = x - \beta t$.

A substitution of the equation (2.6) to equation (2.10) into equation (3.24), leads to the following five traveling wave solutions of the Benjamin-Bona-Mahony equation.



When $\mu \neq 0, \lambda^2 - 4\mu > 0$, the solution of equation (3.16) is

$$\begin{aligned}
 v_6(\xi) = & \frac{\lambda^2 b \beta + 8 \mu b \beta + \beta - 1}{a} \\
 & - \frac{12 b \beta \lambda}{a} \left(\frac{2 \mu}{\sqrt{\lambda^2 - 4 \mu} \tanh\left(\frac{\sqrt{\lambda^2 - 4 \mu}}{2}\right)(\xi + E) + \lambda} \right) \\
 & + \frac{12 b \beta}{a} \left(\frac{2 \mu}{\sqrt{\lambda^2 - 4 \mu} \tanh\left(\frac{\sqrt{\lambda^2 - 4 \mu}}{2}\right)(\xi + E) + \lambda} \right)^2.
 \end{aligned} \tag{3.25}$$

When $\mu \neq 0, \lambda^2 - 4\mu < 0$, the solution of equation (3.16) is

$$\begin{aligned}
 v_7(\xi) = & \frac{\lambda^2 b \beta + 8 \mu b \beta + \beta - 1}{a} \\
 & + \frac{12 b \beta \lambda}{a} \left(\frac{2 \mu}{\sqrt{4 \mu - \lambda^2} \tan\left(\frac{\sqrt{4 \mu - \lambda^2}}{2}\right)(\xi + E) - \lambda} \right) \\
 & + \frac{12 b \beta}{a} \left(\frac{2 \mu}{\sqrt{4 \mu - \lambda^2} \tan\left(\frac{\sqrt{4 \mu - \lambda^2}}{2}\right)(\xi + E) - \lambda} \right)^2.
 \end{aligned} \tag{3.26}$$

When $\mu = 0, \lambda \neq 0, \lambda^2 - 4\mu > 0$, the solution of equation (3.16) is

$$\begin{aligned}
 v_8(\xi) = & \frac{\lambda^2 b \beta + 8 \mu b \beta + \beta - 1}{a} + \frac{12 b \beta \lambda}{a} \left(\frac{\lambda}{\exp(\lambda(\xi + E)) - 1} \right) \\
 & + \frac{12 b \beta}{a} \left(\frac{\lambda}{\exp(\lambda(\xi + E)) - 1} \right)^2.
 \end{aligned} \tag{3.27}$$

When $\mu \neq 0, \lambda \neq 0, \lambda^2 - 4\mu = 0$, the solution of equation (3.16) is

$$\begin{aligned}
 v_9(\xi) = & \frac{\lambda^2 b \beta + 8 \mu b \beta + \beta - 1}{a} - \frac{12 b \beta \lambda}{a} \left(\frac{\lambda^2(\xi + E)}{2(\lambda(\xi + E)) + 2} \right) \\
 & + \frac{12 b \beta}{a} \left(\frac{\lambda^2(\xi + E)}{2(\lambda(\xi + E)) + 2} \right)^2.
 \end{aligned} \tag{3.28}$$

When $\mu = 0, \lambda = 0, \lambda^2 - 4\mu = 0$, the solution of equation (3.16) is

$$v_{10}(\xi) = \frac{\lambda^2 b \beta + 8 \mu b \beta + \beta - 1}{a} + \frac{12 b \beta \lambda}{a} \left(\frac{1}{(\xi + E)} + \frac{12 b \beta}{a} \frac{1}{(\xi + E)^2} \right). \tag{3.29}$$

4. GRAPHICAL REPRESENTATIONS AND NUMERICAL EXPERIMENTS

In this section, we will provide the physical meaning of the obtained solutions of the equation (3.1) and equation (3.16) by using the mathematical software Maple, which are represented in Figure 1 to Figure 10. For further understanding the physical meaning of the solutions, we depicted 3D plot of visualize phenomena. We observe that in the numerical experimentations of the obtained solutions, Figure 1 is bright breather solutions, Figure 2 is the dark breather solutions, Figure 3 is bright rogue waves solutions, Figure 4 is also bright breather solutions, Figure 5 is multiple soliton solutions, Figure 6 is also bright breather solutions, Figure 7 is also multiple



FIGURE 1. Bright breathers solution $v_1(\xi)$ of equation (3.1) with $\alpha = -1$, $\beta = -1$, $\mu = 1$, $\lambda = 3$, $E = 1$ and $A_0 = 1$.

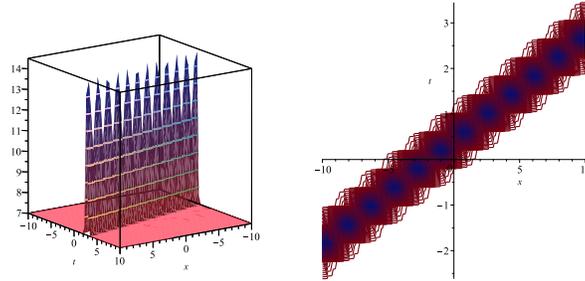


FIGURE 2. Dark breathers solution $v_2(\xi)$ of equation (3.1) with $\alpha = -5$, $\beta = -2$, $\mu = 3$, $\lambda = 1$, $E = 1$ and $A_0 = 5$.

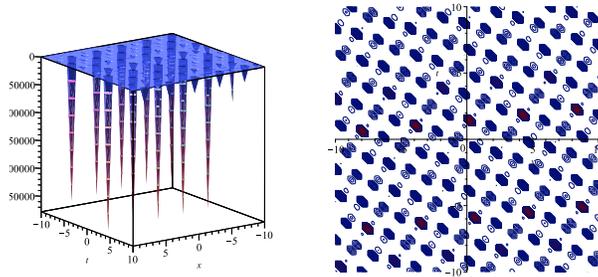
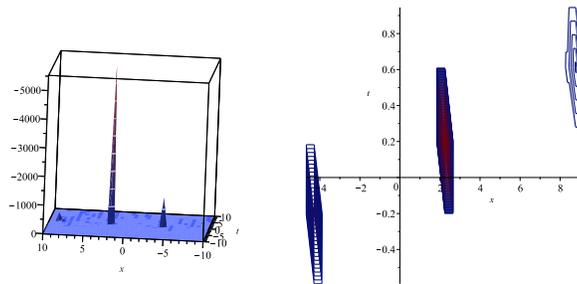


FIGURE 3. Bright rogue waves solution $v_3(\xi)$ of equation (3.1) with $\alpha = -9$, $\beta = -8$, $\mu = 0$, $\lambda = 2$, $E = 1$ and $A_0 = 12$.



soliton solutions, Figure 8 is cuspon waves solitons, Figure 9 is Singular kink waves solutions, and Figure 10 is soliton solutions.



FIGURE 4. Bright breathers solution $v_4(\xi)$ of equation (3.1) with $\alpha = -3, \beta = -3, \mu = 1, \lambda = 2, E = 5$ and $A_0 = 4$.

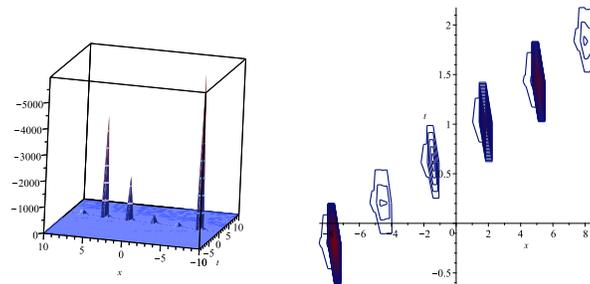


FIGURE 5. Multiple Soliton solution $v_5(\xi)$ of equation (3.1) with $\alpha = -8, \beta = -9, \mu = 0, \lambda = 0, E = 1$ and $A_0 = 4$.

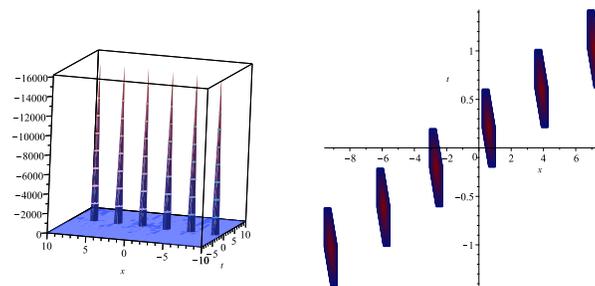
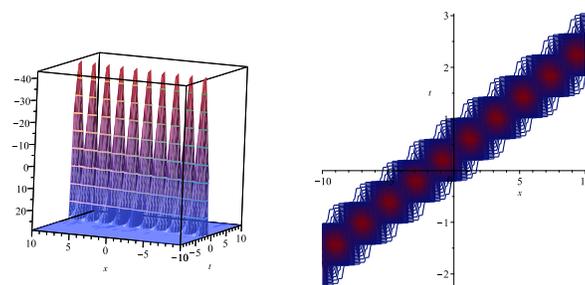


FIGURE 6. Bright breathers solution $v_6(\xi)$ of equation (3.16) with $a = 1, b = 1, \mu = 1, \lambda = 3, E = 1$ and $\beta = 5$.



5. CONCLUSION

In summary, applying the $exp(-\phi(\xi))$ -expansion method to the fourth order Benjamin-Ono equation and BBM equation, we have successfully obtained a special kinds of



FIGURE 7. Multiple bright Soliton solution $v_7(\xi)$ of equation (3.16) with $a = 1$, $b = 1$, $\mu = 3$, $\lambda = 1$, $E = 1$ and $\beta = 1$.

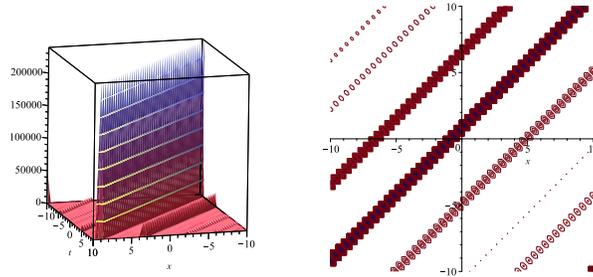


FIGURE 8. Cusp waves solution $v_8(\xi)$ of equation (3.16) with $a = 2$, $b = 2$, $\mu = 0$, $\lambda = 2$, $E = 1$ and $\beta = 1$.

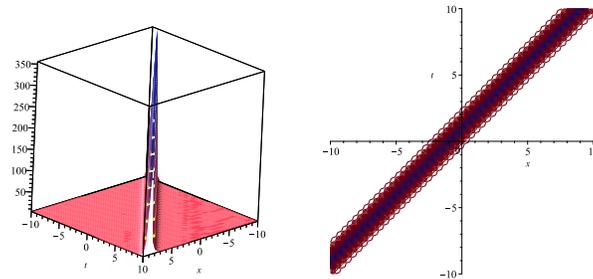
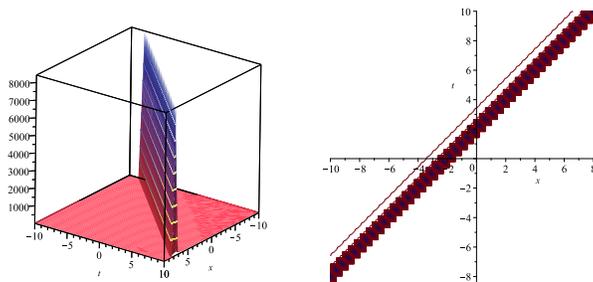


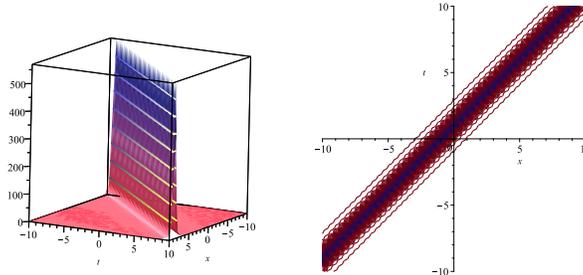
FIGURE 9. Singular kink waves solution $v_9(\xi)$ of equation (3.16) with $a = 1$, $b = 1$, $\mu = 1$, $\lambda = 2$, $E = 1$ and $\beta = 1$.



waves solutions such as bright bright breathers waves, dark breathers waves, a cusp wave, kink waves, bright rogue waves and some soliton solutions.



FIGURE 10. Soliton solution $v_{10}(\xi)$ of equation (3.16) with $a = 5$, $b = 8$, $\mu = 0$, $\lambda = 0$, $E = 1$ and $\beta = 1$.



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